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Crystal, Molecular and Electronic Structure of 1,1-Diaryl-2-halogenoethylenes. I. 2-Bromo-1,1-di-p-tolylethylene

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The crystal structure of 2-bromo-1,1-di-p-tolylethylene, $C_{16}H_{15}Br$, at room temperature, has been determined by three-dimensional methods, including anisotropic refinement by full-matrix least-squares calculations (final R value, 9.7%). There are four formula units in the orthorhombic cell, $a=16.89$, $b=13.07$, $c=6.26$ Å, with space group $P2_12_12_1$. In the molecule the atoms are distributed on three planes: the ethylene reference plane and the two tolyl rings which are rotated by 24.4° (*trans* to Br) and 67.9° (*cis* to Br). Bond lengths and angles are in the normal range; the C-Br distance is 1.87 Å. Bond angles and distances in the isolated molecule and in 'model' crystals were calculated by semi-theoretical methods and the correlation with experimental values is discussed.

Introduction

It is known that 1,1-diaryl-2-halogenoethylenes can react with sodium ethoxide following two different paths leading to nucleophilic substitution and to a Fritsch rearrangement (Beltrame & Favini, 1963; Beltrame, Carrà, Macchi & Simonetta, 1964). For the latter reaction, a geometrical and electronic model of the activated complex has been suggested (Simonetta & Carrà, 1963) in which the aryl group *trans* to the halogen lies in a plane perpendicular to the ethylenic bond. Since the geometry of the substrate might be a relevant factor in determining the relative amounts of the two possible reactions with the base, an investigation was undertaken of the crystal, molecular and electronic structure of some diarylethylenes, including 2-bromo-1,1-di-p-tolylethylene, 2-chloro-1,1-di-p-tolylethylene, 1,1-diphenylethylene and 2-bromo-1,1-di-phenylprop-1-ene.

In this paper the results for 2-bromo-1,1-di-p-tolylethylene are given. The geometry of the molecule, as determined by X-ray diffraction, has been compared with the geometry of the isolated molecule calculated by semi-empirical methods and the interrelations between the two have been discussed.

Experimental

Crystal data

2-Bromo-1,1-di-p-tolylethylene, m.p. 52°C, crystallizes from isopropyl alcohol in transparent plates. The unit-cell dimensions were determined from zero-layer Weissenberg films about the b and c axes, with Cu $K\alpha$ radiation. The crystal data given below have already been published in a preliminary communication (Mariani, Mugnoli & Casalone, 1965).

2-Bromo-1,1-di-p-tolylethylene

$C_{16}H_{15}Br$ F.W. 287.2

Orthorhombic, $a=16.89 \pm 0.01$, $b=13.07 \pm 0.01$, $c=6.26 \pm 0.01$ Å, with $\lambda(\text{Cu } K\alpha_1)=1.5405$, $\lambda(\text{Cu } K\alpha_2)=1.5443$, $\lambda(\text{Cu } K\alpha)=1.5418$ Å; $V=1381.9$ Å³; $D_m=1.37$ g.cm⁻³ (by flotation); $Z=4$; $D_c=1.38$ g.cm⁻³; $F(00)=584$; Absorption coefficient for Cu $K\alpha$ radiation, 39.1 cm⁻¹. Space group $P2_12_12_1$ (no. 19), from systematic absences: $h00$ if $h=2n+1$, $0k0$ if $k=2n+1$, $00l$ if $l=2n+1$.

Intensity measurements

The X-ray intensities were estimated visually from sets of multiple-film equi-inclination Weissenberg

photographs taken at room temperature about the *c* axis (6 layers) and the *b* axis (4 layers). Two crystals were used, with cross sections respectively of $0.16 \times 0.34 \text{ mm}^2$ and of $0.15 \times 0.15 \text{ mm}^2$. From a total number of 1835 possible reflexions, with Cu $K\alpha$ radiation, 1705 were collected, of which 505 were too low to be observed. The intensities were scaled together within the same layer by a method similar to the one described by Rae (1965) and then corrected for Lorentz, polarization and spot-size (Phillips, 1954) factors. An evaluation of the σ for each individual observation was obtained from a statistical analysis (Gramaccioli & Mariani, 1965). The data belonging to different layers were scaled together according to the method proposed by Rollett & Sparks (1960), modified in order to account for the dependence of the weights upon the final scale factors (Hamilton, Rollett & Sparks, 1965; Duchamp, 1964). 242 reflexions were measured about both axes. No absorption or extinction corrections were applied.

Atomic scattering factors were calculated by means of analytical expressions following Cromer & Waber (1965) for bromine and carbon, and Forsyth & Wells (1959) for hydrogen atoms.

All the calculations except the least-squares refinement were carried out on an IBM 1620-20K electronic computer. Structure factors and Fourier syntheses were

calculated by the programs of Albano, Bellon & Pompa (1963) and Pompa, Albano, Bellon & Scatturin (1963). The intensity scaling program was written by Gramaccioli and Mariani, who also adapted and partially modified the least-squares ORFLS program of Busing, Martin & Levy (1962) for an IBM 7040 computer. All the programs for the theoretical calculations were written by one of us (C.M.).

Structure determination and refinement

A preliminary two-dimensional investigation by Patterson and Fourier methods on the (*hk0*) and the (*h0l*) projections (Mariani, Mugnoli & Casalone, 1965) gave a set of atomic coordinates and isotropic thermal factors, used as a starting point for the refinement. After two cycles of differential Fourier synthesis, and one three-dimensional difference synthesis calculation, final refinement was achieved by some cycles of full-matrix least-squares, using the program ORFLS written by Busing, Martin & Levy (1962). During these last cycles, anisotropic temperature factors were adopted for each atom in the structure, except for hydrogen atoms. These latter were omitted from the least-squares refinement, being introduced only in the structure factor calculations. Hydrogen atoms bonded to an sp^2 carbon atom were assumed to lie on the line bisecting the angle

Table 1. Final coordinates with standard deviations ($\times 10^4$)

	<i>x/a</i>	<i>y/b</i>	<i>z/c</i>		<i>x/a</i>	<i>y/b</i>	<i>z/c</i>
Br	1172 (1)	2463 (1)	3738 (2)	C(13)	1942 (6)	5919 (7)	977 (20)
C(1)	464 (7)	2834 (8)	1574 (22)	C(14)	1478 (6)	5104 (9)	254 (19)
C(2)	407 (6)	3790 (7)	799 (16)	C(15)	-1847 (8)	4779 (12)	-5898 (24)
C(3)	-177 (6)	4015 (7)	-941 (15)	C(16)	2341 (9)	7162 (9)	3853 (24)
C(4)	-445 (7)	3266 (7)	-2350 (19)	H(1)	82	2251	915
C(5)	-996 (7)	3518 (8)	-3939 (18)	H(2)	-228	2491	-2218
C(6)	-1261 (6)	4521 (10)	-4162 (22)	H(3)	-1216	2931	-4997
C(7)	-987 (6)	5272 (9)	-2767 (20)	H(4)	-1190	6051	-2933
C(8)	-449 (5)	5019 (7)	-1151 (17)	H(5)	-243	5602	-62
C(9)	913 (6)	4674 (7)	1602 (19)	H(6)	370	4702	4710
C(10)	808 (7)	5042 (9)	3672 (18)	H(7)	1196	6146	5994
C(11)	1274 (8)	5855 (8)	4393 (22)	H(8)	2378	6259	-65
C(12)	1840 (7)	6292 (7)	3049 (21)	H(9)	1557	4809	-1343

Table 2. Thermal exponent coefficients (with e.s.d.'s)

b_{ij} as given here are defined by $T = \exp [-10^{-4}(b_{11}h^2 + b_{22}k^2 + b_{33}l^2 + 2b_{12}hk + 2b_{13}hl + 2b_{23}kl)]$

	b_{11}	b_{22}	b_{33}	b_{12}	b_{13}	b_{23}
Br	76 (1)	75 (1)	651 (6)	-3 (1)	-5 (2)	71 (2)
C(1)	55 (4)	65 (6)	571 (40)	-10 (4)	4 (12)	25 (14)
C(2)	48 (4)	47 (4)	419 (29)	-3 (3)	11 (9)	0 (10)
C(3)	40 (3)	62 (5)	347 (28)	-3 (4)	9 (8)	-15 (10)
C(4)	59 (5)	57 (6)	486 (34)	-9 (5)	30 (12)	-10 (12)
C(5)	58 (5)	71 (6)	503 (37)	-21 (5)	11 (12)	-34 (14)
C(6)	41 (4)	90 (7)	426 (30)	-6 (5)	-21 (9)	-27 (14)
C(7)	36 (4)	75 (6)	484 (32)	6 (4)	-2 (9)	-49 (12)
C(8)	36 (3)	62 (5)	489 (33)	12 (3)	-6 (9)	-70 (12)
C(9)	39 (3)	51 (5)	419 (30)	-6 (3)	8 (8)	22 (10)
C(10)	44 (4)	84 (7)	418 (33)	-1 (4)	10 (9)	45 (14)
C(11)	72 (6)	51 (5)	456 (29)	10 (5)	-28 (12)	3 (11)
C(12)	57 (5)	48 (5)	565 (41)	8 (4)	-67 (12)	-11 (12)
C(13)	41 (4)	50 (5)	539 (37)	-7 (3)	-11 (10)	21 (12)
C(14)	40 (3)	68 (7)	525 (36)	6 (4)	25 (10)	-37 (13)
C(15)	69 (6)	125 (11)	459 (38)	-22 (7)	-13 (14)	-4 (18)
C(16)	77 (6)	65 (6)	646 (49)	-9 (5)	-49 (16)	-19 (16)

STRUCTURE OF 1,1-DIARYL-2-HALOGENOETHYLENES. I

Table 3. Observed and calculated structure factors

Each group of four columns contains in order h , $10|F_0|$, $10|F_c|$, and the phase angle. A minus sign before the $10|F_0|$ value designates unobserved reflexion. Reflexions indicated by an asterisk were given zero weight in the least-squares refinement and excluded from the R calculation.

H	0 0	12 227 224 0	13 24 17 90	6 257 294 133	13 49 63 82	0 115 2 22	4 103 95 289	6 230 221 270	13 87 99 80			
* 2 623 782 0	14 38 33 90	14 257 294 133	7 146 119 124	14 64 61 251	5 214 250 191	7 30 27 270	14 -19 12 102					
4 1215 1442 180	15 38 47 270	8 257 295 86	15 -21 22 154	* 1 661 752 185	6 70 67 79	8 21 23 90	15 -16 12 276					
5 1392 1600 180	15 38 47 270	0 135 129 180	16 237 294 180	16 46 41 214	7 243 143 180	7 64 63 123	9 103 102 90	16 21 25 92				
6 254 241 0	17 -33 180	10 237 294 180	11 237 295 275	17 -13 24 264	2 157 141 80	9 -35 26 166	10 121 120 90	17 20 25 92				
10 251 241 0	18 -30 34 180	2 43 47 180	10 133 124 180	10 133 124 180	11 161 143 80	11 -38 26 166	12 24 10 90	18 -11 15 271				
12 392 363 180	19 -24 23 90	3 70 64 270	13 54 5 240	14 265 300 90	5 248 250 35	11 123 129 352	13 67 82 270					
14 201 251 180	19 -24 23 90	4 146 134 270	14 103 118 115	14 103 118 115	6 263 300 90	12 123 129 352	14 57 82 270	0 265 314 270				
15 46 0	19 -24 23 90	5 146 134 270	15 237 294 180	2 97 57 264	7 263 154 90	13 -33 31 20	15 24 27 270	1 265 228 141				
18 -39 4 180	20 27 0	17 -33 17 0	16 237 294 180	16 237 295 113	9 111 104 204	14 -33 31 20	16 -19 14 270	2 89 77 217				
20 27 20 0	1 27 28 90	7 33 43 90	17 237 32 81	12 154 286	10 175 151 150	15 30 34 187	17 -16 9 90	2 243 229 59				
1 270 324 270	2 164 599 59	8 121 185 180	18 -24 8 36	8 124 130 231	16 27 247 193	18 30 32 90	19 -30 1 180	2 275 290 68				
2 750 785 180	2 164 265 215 90	9 121 185 180	19 -24 8 36	9 124 130 231	17 27 247 193	19 -30 1 180	20 1 180	2 275 290 68				
3 475 571 180	4 164 473 429 180	12 -13 18 0	10 113 111 233	10 113 111 233	11 123 129 266	12 123 129 266	13 123 129 266	7 49 45 213				
4 168 136 180	4 164 473 429 180	12 -13 18 0	9 -36 111 106	12 123 129 266	13 123 129 266	14 123 129 266	15 123 129 266	8 146 153 209				
5 151 175 90	7 -35 111 180	11 15 0	0 634 562 270	11 15 0	1 214 188 335	1 227 232 276	1 227 232 276	9 91 84 241				
6 912 821 0	9 30 27 270	1 46 40 270	1 338 357 269	12 49 47 268	18 -30 16 146	2 235 223 3	1 91 83 94					
7 -24 13 270	10 292 256 0	2 167 147 0	2 -19 15 348	13 -24 18 196	19 -27 12 146	3 232 223 196	12 61 59 126					
8 146 144 0	11 -34 28 90	5 51 38 270	14 363 392 288	15 51 45 242	20 -21 11 1369	4 232 223 196	13 -19 29 27					
9 146 144 0	11 -34 28 90	5 51 38 270	15 51 45 242	21 -11 1369	5 278 262 159	6 64 65 233	6 87 268 266	15 31 35 258				
10 359 342 181	13 -34 21 70	5 43 16 55	16 -38 14 355	17 -13 10 33	7 121 134 266	8 181 176 36	8 214 250 71	17 -11 9 249				
11 166 164 91	14 158 153 180	6 94 86 180	15 160 153 180	16 127 139 260	7 181 176 36	8 214 250 71	17 -11 9 249					
12 114 199 180	15 160 153 180	6 94 86 180	16 -24 2 90	17 237 294 180	7 181 176 36	8 214 250 71	17 -11 9 249					
13 125 150 90	16 150 153 180	6 94 86 180	17 237 294 180	18 237 294 180	7 181 176 36	8 214 250 71	17 -11 9 249					
14 125 150 90	16 46 27 21 0	10 111 53	11 111 53	12 111 53	13 123 129 266	14 123 129 266	15 123 129 266	16 123 129 266				
15 64 75 270	16 -27 21 0	10 111 53	11 111 53	12 111 53	13 123 129 266	14 123 129 266	15 123 129 266	16 123 129 266				
16 113 98 0	19 -21 20 270	1 143 131 270	2 -24 21 0	3 114 140 150	4 173 165 251	5 27 34 175	6 35 39 55	1 114 126 19				
20 -27 25 180	2 136 107 0	3 114 140 150	4 173 165 251	5 27 34 175	6 35 39 55	7 178 176 36	8 214 250 71	9 175 183 269				
21 -21 11 90	3 114 140 150	4 173 165 251	5 27 34 175	6 35 39 55	7 178 176 36	8 214 250 71	9 175 183 269	10 178 183 269				
H	2 0	11 8 0	0 0	0 113 111 0	1 12 1 1	0 278 264 90	0 278 264 90	0 278 264 90	H 8 3			
1 205 266 180	2 151 126 90	3 116 20 270	4 -16 2 90	5 127 326 117	6 127 326 117	7 127 326 117	8 127 326 117	9 116 112 65				
1 626 513 270	6 103 87 180	7 -11 20 270	8 127 326 117	9 127 326 117	10 127 326 117	11 127 326 117	12 127 326 117	13 116 112 65				
3 781 785 270	8 127 326 117	9 127 326 117	10 127 326 117	11 127 326 117	12 127 326 117	13 127 326 117	14 127 326 117	15 116 112 65				
4 257 307 9	9 111 20 270	* 1 389 327 270	0 113 130 0	1 130 130 0	1 130 130 0	1 130 130 0	1 130 130 0	1 130 130 0	H 8 3			
5 130 144 90	10 111 93 0	* 2 286 327 270	1 499 534 286	2 610 624 285	3 728 737 270	4 73 73 160	5 74 74 160	6 75 75 160	7 116 112 65			
6 332 270 0	11 -38 3 270	5 738 877 270	6 281 266 98	0 118 163 270	1 124 127 260	2 125 127 260	3 126 127 260	4 127 127 260	5 116 112 65			
7 -24 9 270	9 120 108 180	10 113 180 180	11 113 180 180	12 113 180 180	13 113 180 180	14 113 180 180	15 113 180 180	16 113 180 180	17 113 180 180			
8 632 595 180	13 -38 4 270	14 108 180 180	15 108 180 180	16 108 180 180	17 108 180 180	18 108 180 180	19 108 180 180	20 108 180 180	21 108 180 180			
9 146 146 90	14 -38 4 270	15 108 180 180	16 108 180 180	17 108 180 180	18 108 180 180	19 108 180 180	20 108 180 180	21 108 180 180	22 108 180 180			
10 91 72 180	15 -35 3 270	16 108 180 180	17 -24 2 90	18 108 180 180	19 108 180 180	20 108 180 180	21 108 180 180	22 108 180 180	23 108 180 180			
11 91 72 180	16 -35 3 270	17 -24 2 90	18 108 180 180	19 108 180 180	20 108 180 180	21 108 180 180	22 108 180 180	23 108 180 180	24 108 180 180			
12 91 72 180	17 -24 2 90	18 108 180 180	19 108 180 180	20 108 180 180	21 108 180 180	22 108 180 180	23 108 180 180	24 108 180 180	25 108 180 180			
13 91 72 180	18 -24 2 90	19 108 180 180	20 108 180 180	21 108 180 180	22 108 180 180	23 108 180 180	24 108 180 180	25 108 180 180	26 108 180 180			
14 91 72 180	19 -24 2 90	20 108 180 180	21 108 180 180	22 108 180 180	23 108 180 180	24 108 180 180	25 108 180 180	26 108 180 180	27 108 180 180			
15 91 72 180	20 108 180 180	21 108 180 180	22 108 180 180	23 108 180 180	24 108 180 180	25 108 180 180	26 108 180 180	27 108 180 180	28 108 180 180			
16 91 72 180	21 108 180 180	22 108 180 180	23 108 180 180	24 108 180 180	25 108 180 180	26 108 180 180	27 108 180 180	28 108 180 180	29 108 180 180			
17 91 72 180	22 108 180 180	23 108 180 180	24 108 180 180	25 108 180 180	26 108 180 180	27 108 180 180	28 108 180 180	29 108 180 180	30 108 180 180			
18 91 72 180	23 108 180 180	24 108 180 180	25 108 180 180	26 108 180 180	27 108 180 180	28 108 180 180	29 108 180 180	30 108 180 180	31 108 180 180			
19 91 72 180	24 108 180 180	25 108 180 180	26 108 180 180	27 108 180 180	28 108 180 180	29 108 180 180	30 108 180 180	31 108 180 180	32 108 180 180			
20 91 72 180	25 108 180 180	26 108 180 180	27 108 180 180	28 108 180 180	29 108 180 180	30 108 180 180	31 108 180 180	32 108 180 180	33 108 180 180			
21 91 72 180	26 108 180 180	27 108 180 180	28 108 180 180	29 108 180 180	30 108 180 180	31 108 180 180	32 108 180 180	33 108 180 180	34 108 180 180			
H	3 0	11 8 0	0 0	0 113 111 0	0 114 1 1	0 115 2 2	0 116 3 2	0 117 4 2	0 118 5 2	0 119 6 2		
* 1 238 159 90	7 -35 143 0	8 148 143 0	9 148 143 0	10 148 143 0	11 148 143 0	12 148 143 0	13 148 143 0	14 148 143 0	15 148 143 0	16 148 143 0		
3 543 432 270	10 230 229 180	H 1 1	H 1 1	H 1 1	H 1 1	H 1 1	H 1 1	H 1 1	H 1 1	H 1 1		
4 126 164 180	11 10 0	0 426 456 270	0 227 247 90	1 191 205 244	2 243 252 290	3 243 252 290	4 243 252 290	5 243 252 290	6 243 252 290	7 243 252 290		
5 154 130 90	12 -14 116 180	1 372 450 252	1 262 450 252	2 262 450 252	3 262 450 252	4 262 450 252	5 262 450 252	6 262 450 252	7 262 450 252	8 262 450 252		
6 824 772 180	13 -34 166 180	* 2 307 408 8	2 241 241 348	3 241 241 348	4 241 241 348	5 241 241 348	6 241 241 348	7 241 241 348	8 241 241 348	9 241 241 348		
7 824 285 90	14 -35 148 48 0	* 2 307 408 8	2 241 241 348	3 241 241 348	4 241 241 348	5 241 241 348	6 241 241 348	7 241 241 348	8 241 241 348	9 241 241 348		
8 916 342 180	15 -35 148 48 0	* 2 307 408 8	2 241 241 348	3 241 241 348	4 241 241 348	5 241 241 348	6 241 241 348	7 241 241 348	8 241 241 348	9 241 241 348		
9 517 468 0	16 -35 148 48 0	* 2 307 408 8	2 241 241 348	3 241 241 348	4 241 241 348	5 241 241 348	6 241 241 348	7 241 241 348	8 241 241 348	9 241 241 348		
10 465 398 0	17 -35 148 48 0	* 2 307 408 8	2 241 241 348	3 241 241 348	4 241 241 348	5 241 241 348	6 241 241 348	7 241 241 348	8 241 241 348	9 241 241 348		
11 91 75 90	18 -38 148 0	* 2 307 408 8	2 241 241 348	3 241 241 348	4 241 241 348	5 241 241 348	6 241 241 348	7 241 241 348	8 241 241 348	9 241 241 348		
12 259 223 180	19 -38 148 0	* 2 307 408 8	2 241 241 348	3 241 241 348	4 241 241 348	5 241 241 348	6 241 241 348	7 241 241 348	8 241 241 348	9 241 241 348		
13 259 223 180	20 -17 20 270	1 318 351 89	2 316 351 89									

Table 3 (cont.)

-8	25	227	16	-19	19	39	5	54	50	51	342	1	27	36	138	7	-19	9	90	0	H	4	5	6	180	3	-33	18	90								
1	15	3	16	-19	17	32	5	7	27	29	34	2	22	31	162	7	-19	9	90	0	H	4	5	6	180	3	-33	18	90								
0	1	44	27	11	-21	26	5	7	151	153	154	3	165	179	181	10	-19	81	91	95	0	H	4	5	6	180	3	-33	18	90							
1	73	88	127	1	-21	30	91	4	-21	29	252	5	-21	14	81	11	-19	21	270	3	H	4	5	6	180	3	-33	18	90								
2	-11	3	245	1	31	34	101	61	82	297	6	116	113	111	11	-19	19	270	3	H	4	5	6	180	3	-33	18	90									
3	63	61	43	2	32	35	178	12	87	14	54	7	121	117	20	14	-19	27	90	3	H	4	5	6	180	3	-33	18	90								
4	-7	47	317	1	33	56	59	93	13	-21	26	32	9	-19	31	255	15	-16	9	90	0	H	4	5	6	180	3	-33	18	90							
5	6	-13	64	4	73	73	142	14	21	26	32	10	-13	42	43	337	16	-13	9	90	0	H	4	5	6	180	3	-33	18	90							
6	-6	1	100	119	134	187	201	10	-13	23	11	1	10	25	174	17	-13	8	228	0	H	4	5	6	180	3	-33	18	90								
7	0	5	4	575	55	178	180	7	135	141	164	262	17	-11	9	304	14	-6	12	176	0	H	4	5	6	180	3	-33	18	90							
8	1	38	47	180	7	87	92	91	110	107	265	1	167	163	270	0	11	4	5	27	1	H	4	5	6	180	3	-33	18	90							
9	2	70	74	88	184	0	9	114	107	265	10	135	125	165	0	15	27	268	0	H	4	5	6	180	3	-33	18	90									
10	3	57	67	72	0	110	101	107	141	147	164	2	208	212	191	1	30	28	58	1	H	4	5	6	180	3	-33	18	90								
11	5	24	248	1	12	37	81	167	3	46	55	81	2	127	116	79	5	-19	121	141	261	0	H	4	5	6	180	3	-33	18	90						
12	97	11	176	180	1	13	54	53	73	3	46	55	81	3	-21	32	245	6	99	96	313	0	H	4	5	6	180	3	-33	18	90						
13	7	111	144	180	1	14	46	63	321	4	91	94	113	5	-19	121	141	261	6	H	4	5	6	180	3	-33	18	90									
14	8	124	180	1	15	34	36	101	54	55	81	6	18	16	79	5	-19	121	141	261	0	H	4	5	6	180	3	-33	18	90							
15	9	-2	54	58	180	1	16	13	115	16	43	45	253	6	-19	18	166	2	-19	17	47	0	H	4	5	6	180	3	-33	18	90						
16	11	-21	5	0	18	11	-21	22	26	8	79	89	329	7	-19	26	268	10	7	43	44	111	111	H	4	5	6	180	3	-33	18	90					
17	12	63	65	1	12	21	38	101	30	31	198	8	-21	39	338	11	-19	34	46	10	H	4	5	6	180	3	-33	18	90								
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22	17	1	111	119	169	360	7	81	67	29	167	1	24	18	67	0	11	21	25	180	12	H	4	5	6	180	3	-33	18	90							
23	18	33	39	180	8	81	121	102	149	0	231	240	180	3	-19	22	30	13	1	H	4	5	6	180	3	-33	18	90									
24	19	26	180	9	-21	22	24	101	54	55	81	9	165	165	165	1	-19	19	23	19	1	H	4	5	6	180	3	-33	18	90							
25	20	3	49	40	85	11	-21	33	328	3	110	121	256	6	-21	32	304	3	54	56	92	0	H	4	5	6	180	3	-33	18	90						
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30	25	103	101	252	17	13	-16	38	16	4	56	53	124	11	-19	21	28	7	21	30	93	6	H	4	5	6	180	3	-33	18	90						
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37	32	7	6	7	5	3	73	73	73	23	-13	72	178	5	-16	16	54	141	17	13	20	239	13	H	4	5	6	180	3	-33	18	90					
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39	34	7	10	76	6	202	227	191	0	71	9	8	90	9	-11	29	116	15	2	97	98	245	0	H	4	5	6	180	3	-33	18	90					
40	35	7	10	76	6	202	227	191	0	71	9	8	90	9	-11	29	116	15	2	97	98	245	0	H	4	5	6	180	3	-33	18	90					
41	36	7	10	76	6	202	227	191	0	71	9	8	90	9	-11	29	116	15	2	97	98	245	0	H	4	5	6	180	3	-33	18	90					
42	37	0	1	40	38	209	24	34	46	50	214	0	49	45	40	78	70	141	2	81	81	149	0	H	4	5	6	180	3	-33	18	90					
43	38	1	10	81	87	209	34	40	52	282	0	49	45	40	78	70	141	2	81	81	149	0	H	4	5	6	180	3	-33	18	90						
44	39	1	10	81	87	209	34	40	52	282	0	49	45	40	78	70	141	2	81	81	149	0	H	4	5	6	180	3	-33	18	90						
45	40	2	10	81	87	209	34	40	52	282	0	49	45	40	78	70	141	2	81	81	149	0	H	4	5	6	180	3	-33	18	90						
46	41	2	10	81	87	209	34	40	52	282	0	49	45	40	78	70	141	2	81	81	149	0	H	4	5	6	180	3	-33	18	90						
47	42	2	10	81	87	209	34	40	52	282	0	49	45	40	78	70	141	2	81	81	149	0	H	4	5	6	180	3	-33	18	90						
48	43	2	10	81	87	209	34	40	52	282	0	49	45	40	78	70	141	2	81	81	149	0	H	4	5	6	180	3	-33	18	90						
49	44	2	10	81	87	209	34	40	52	282	0	49	45	40	78	70	141	2	81	81	149	0	H	4	5	6	180	3	-33	18	90						
50	45	2	10	81	87	209	34	40	52	282	0	49	45	40	78	70	141	2	81	81	149	0	H	4	5	6	180	3	-33	18	90						
51	46	2	10	81	87	209	34	40	52	282	0	49	45	40	78	70	141	2	81	81	149	0	H	4	5	6	180	3	-33								

between the two bonds from the same carbon atom, with carbon-hydrogen bond distances equal to 1.08 Å. The positions of the hydrogen atoms of the methyl groups were not determined.

The quantity minimized in the least-squares calculation was $\sum w(F_o - |kF_c|)^2$, the weights w being derived from the σ 's of each single observation, as determined in the data reduction process. After the final cycle, all the shifts were below the corresponding e.s.d. Unobserved reflexions were assigned a threshold value based on the lowest observable intensity, and included in the refinement only when $|kF_c|$ exceeded this value. The final R index is 9.7% for the observed reflexions.

Table 1 gives the final atomic coordinates with their e.s.d.'s. Anisotropic temperature factors with their e.s.d.'s are listed in Table 2. An isotropic temperature factor $B=6 \text{ \AA}^2$ was assigned to each hydrogen atom. The numbering of the atoms is included in Figs. 1 and 2. Structure factors based on the final parameters are compared with the observed structure amplitudes in Table 3. The average planes in the molecule were determined by the method of Schomaker, Waser, Marsh & Bergman (1959). The equations for the two tolyl planes, including the C(2) atom, are

$$0.7555X + 0.2006Y - 0.6237Z - 1.1981 = 0$$

and

$$0.6820X - 0.6500Y + 0.3352Z + 2.5835 = 0$$

for the ring *trans* and *cis* to the bromine atom respectively; X , Y , Z are coordinates in Å. The equation of

the plane of the molecular skeleton C(1), C(2), C(3), C(6), C(15), C(9), C(12), C(16) is

$$0.6904X - 0.2122Y - 0.6916Z + 0.9222 = 0.$$

In all the three cases, no distance exceeding 0.01 Å was found between any atom and the involved plane. The bromine atom is coplanar with the molecular skeleton within 0.012 Å.

Bond distances, bond angles and the averages of their e.s.d.'s (Cruickshank & Robertson, 1953; Darlow, 1960) are reported in Figs. 1 and 2 respectively. All these values are uncorrected for librational effects. In Figs. 1 and 2 the molecule is drawn referred to an orthogonal, right-handed coordinate system $X'Y'Z'$ whose $Y'Z'$ plane is the average plane of the molecular skeleton previously described. The origin is taken at the mass centre of the skeleton. Both the tolyl rings are rotated with respect to the molecular skeleton plane, the rotation angles being $\theta_1 = 24.4^\circ$ and $\theta_2 = 67.9^\circ$ for the ring *trans* and *cis* to the bromine atom respectively.

In Figs. 3 and 4 the packing in the crystal is viewed along the c and b axes respectively.

π -Electron, bending and steric repulsion energy calculations

The π -electron energy for the molecule with $\theta_1 = 5, 15, 25, 35, 45^\circ$ and $\theta_2 = 45, 55, 65, 75, 85^\circ$ was calculated by the Hückel molecular orbital (HMO) method. For the methyl groups the heteroatom model was used,

with $\alpha_{\text{CH}_3} = \alpha_{\text{C}} + 3.3\beta_0$ and $\beta_{\text{C}-\text{CH}_3} = 0.8\beta_0$ (Favini & Simonetta, 1963) where β_0 is the resonance integral for a carbon–carbon bond at a distance of 1.40 Å. As the σ -bond compression energy was neglected in this calculation, the standard value $\beta_0 = -20 \text{ Kcal.mole}^{-1}$ was chosen. The parameters for bromine were $\alpha_{\text{Br}} = \alpha_{\text{C}} + \beta_0$

and $\beta_{\text{C}-\text{Br}} = 0.4\beta_0$ (Carrà, Raimondi & Simonetta, 1966; Streitwieser, 1962).

Calculations were performed for all the 25 θ_1 , θ_2 combinations by the ω technique, with $\omega = 0.625$, the ω value being chosen to give the quickest convergence. At each step of the calculation the bond orders n_{st}

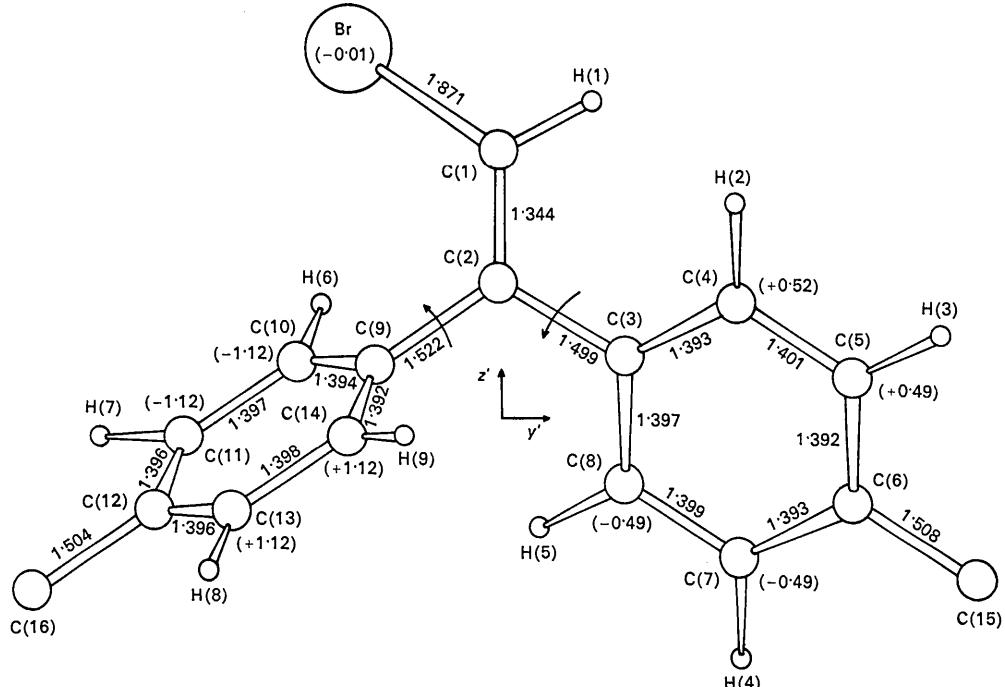


Fig. 1. Bond lengths (\AA) in the crystal. The average e.s.d. is 0.016 \AA . In brackets, the heights (\AA) of non-coplanar atoms are given, with respect to the average plane of the molecular skeleton (see text).

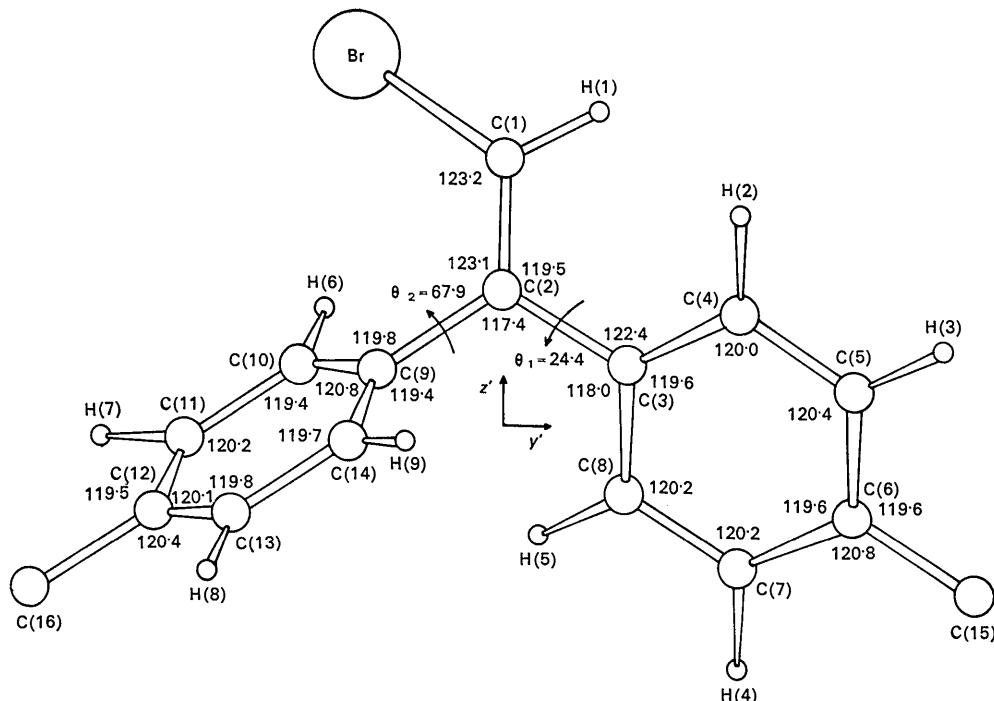


Fig. 2. Bond angles ($^\circ$) in the crystal. The average e.s.d. is 1.0°. Both the tolyl rings are rotated with respect to the average plane of the molecular skeleton (see text).

were evaluated by the standard HMO formula for all bonds except C(2)–C(3) and C(2)–C(9). For these bonds n_{st} was taken as $\sum_i 2c_{si}c_{ti} \cos \theta_j$ ($j=1, 2$; i runs over all the occupied orbitals). The C_{sp^2} –C_{methyl} and C–Br distances were kept at the constant values of 1.51 Å (Keidel & Bauer, 1956) and 1.87 Å. This latter value is based on values found in the literature for vinyl bromide (Whiffen, 1958; Howe & Goldstein, 1957) and 1,2-dibromoethylene (Hugill, Coop & Sutton, 1938). The C_{sp^2} – C_{sp^2} distances were calculated from the bond orders with Coulson's formula (Coulson, 1939), as given by Moser (1953). The integral β_{C-C} was varied with the distance R in the following way

$$\beta(R) = \beta_0 \exp [2.683(1.40 - R)] \quad (1)$$

(Coulson & Golebiewski, 1961) for all the C–C bonds except C(2)–C(3) and C(2)–C(9). For these bonds the

value of β given by (1) was multiplied by $\cos \theta_j$. For each geometry the calculation was repeated to consistency of α 's and β 's in the second decimal, which corresponds to a consistency to the third decimal for distances. As the final values for all C–C distances in the two aromatic rings for all examples examined were in the range 1.395–1.405 Å, the two rings were assumed to be regular hexagons with a side of 1.40 Å.

At this stage four more degrees of freedom were added to our model: each of the angles φ_1 , φ_2 , φ_3 and φ_4 (Fig. 5), which lie in the plane of C(1), C(2), C(3), C(9), Br, and H(1), was allowed to assume one of four values, 115, 120, 125 or 130°. With the 25 possible combinations of torsion angles, this makes the total number of possible geometrical models to be considered 6400. For 972 of them bending and steric repulsion energies were calculated and added to the

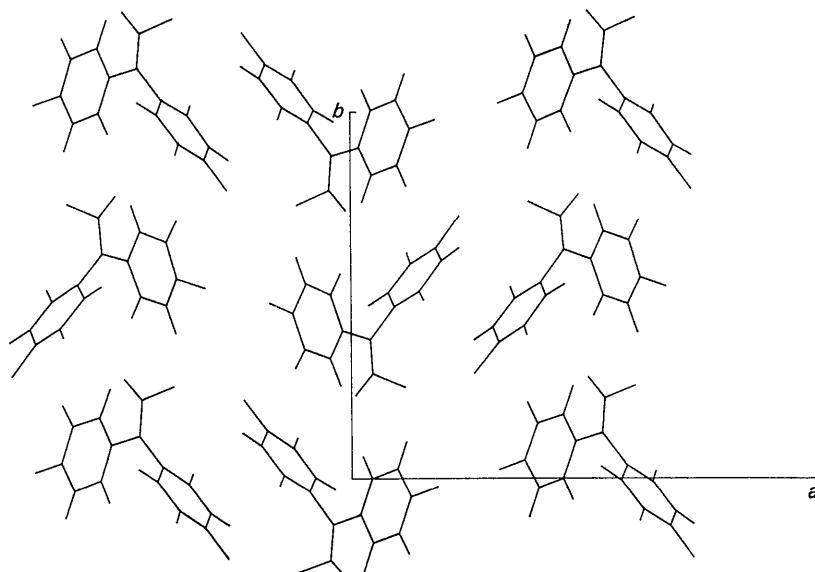


Fig. 3. The crystal structure projected on (001).

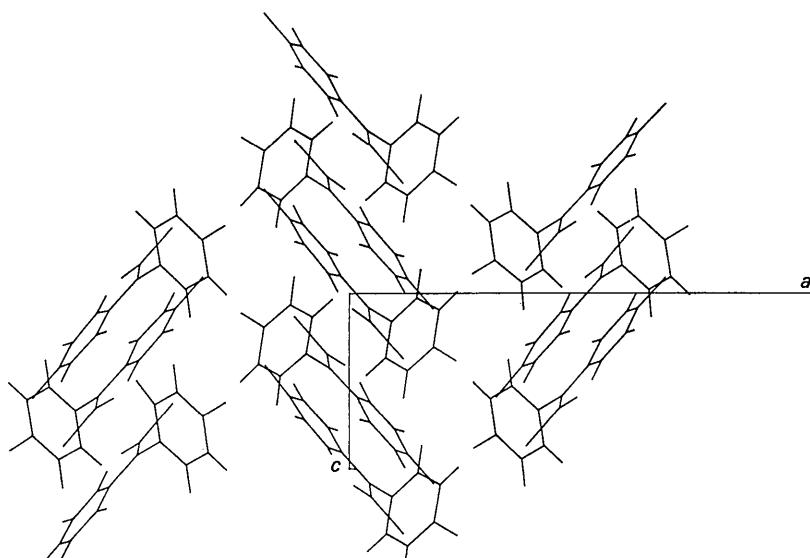


Fig. 4. The crystal structure projected on (010).

appropriate value of the π -electron energy to get a 'total energy' for the isolated molecule. The geometries for which the calculation was actually performed were chosen in such a way as to make possible the identification of the position of the minimum for the function $E=E(\theta_1, \theta_2, \varphi_1, \varphi_2, \varphi_3, \varphi_4)$.

The bending energy in the plane of bonds has been determined by the formula

$$E_{\text{bend}} = \frac{1}{2} \sum_{i=1}^4 k_i (\varphi_i - \varphi_0)^2 \quad (2)$$

where $\varphi_0 = 120^\circ$ and k is the force constant. The following values for the force constants were used: $k_1 = k_2 = 1.35 \cdot 10^{-11}$ erg.rad $^{-2}$ (Favini & Simonetta, 1963), $k_3 = 0.79 \cdot 10^{-11}$ erg.rad $^{-2}$ (Dowling, Puranik, Meister & Miller, 1957), $k_4 = 0.68 \cdot 10^{-11}$ erg.rad $^{-2}$ (Westheimer, 1956). The compression energy for pairs of atoms at distances smaller than the 'touching distance' was calculated by the formula

$$E_s = -0.721(r_0/r)^6 + 2.11 \cdot 10^4 \exp(-10.88r/r_0), \quad (3)$$

where E_s in Kcal.mole $^{-1}$ is the interaction energy, r the actual distance and r_0 the 'touching distance', obtained by adding van der Waals atomic radii. The values of these radii have been taken from standard tables (Pauling, 1960) and corrected when the line joining the two atoms and the direction of the nearest bond make an angle smaller than 90° (Dostrovsky, Hughes & Ingold, 1946). In this calculation the methyl group was considered as a single atom with van der Waals radius equal to 2.0 Å. Equation (3) is an extension of the one given by Kitaygorodsky (1961). This was originally given to calculate carbon and hydrogen interactions and has been here adapted to include methyl groups and bromine atoms. For 12 of the 25 possible sets θ_1 and θ_2 the minimum of energy, with respect to $\varphi_1, \varphi_2, \varphi_3, \varphi_4$, is given in Table 4, the absolute minimum of energy being indicated by an asterisk. The corresponding molecular geometry is also given in Table 4.

However, in the crystal each molecule experiences the strain of the surrounding molecules which causes a distortion of the molecules so that the crystal reaches

Table 4. 'Total energy' (Kcal.mole $^{-1}$)
for the isolated molecule

The most stable conformation, indicated by an asterisk, has bond lengths C(1)-C(2)=1.35, C(2)-C(3)=1.48, C(2)-C(9)=1.52 Å.

θ_1	θ_2	φ_1	φ_2	φ_3	φ_4	Total energy
5	55	125	120	120	125	-674.15
15	45	125	120	125	120	-673.44
15	55	125	120	120	120	-673.83
15	65	125	120	120	120	-673.26
25	45	120	120	125	120	-671.28
25	55	120	120	120	120	-674.18*
25	65	120	120	120	120	-673.27
25	75	120	120	120	120	-673.03
35	45	120	120	125	120	-672.89
35	55	120	120	120	120	-673.39
35	65	120	120	120	120	-673.00
45	55	120	120	120	120	-672.55

the minimum of free energy. We know from X-ray data the position and orientation of the real molecules in the crystal and we put our model molecules in the same places. Considering a reference molecule, we make the C(2) atom, the C(1)-C(2) axis and the C(1), C(2), C(3), C(9) plane of the model coincident with the same atom, axis and plane of a real molecule in the crystal. We then calculate the coordinates of all the atoms in the model molecule, and by symmetry operations we can also get the coordinates of the atoms of all the model molecules in the crystal. We assume that the thermal energy and the entropy of the real crystal and of crystals made from any one of our model molecules are approximately the same. Also the attractive van der Waals forces are about the same, so what we need to do to find the theoretical geometry of the molecules in the crystal is to add to the 'total energy' of the isolated molecule the repulsion energy from the surroundings and then to find the minimum.

We used again formula (3), this time without correction of the van der Waals radii for the angular distance from the nearest bond. This simplification was made necessary by the need of keeping the program of the calculation to a size that the computer could accept and was justified by the fact that, in general, the direction from an atom of the reference molecule to an atom in surrounding molecules makes an angle larger than 90° with the direction of the nearest bond. We made first a trial calculation including all molecules for which C(2) lay within a sphere of radius 20 Å centered on C(2) of the reference molecule. It was recognized that most of the molecules had all the atoms at distances larger than the 'touching distance' with any of the atoms in the reference molecule. It was seen that the appropriate radius for the sphere to include all the molecules between which repulsion may exist was 10 Å, and ten surrounding molecules were partially inside this sphere. This calculation however could not be performed for all the geometries we had considered before. We kept $\varphi_1 = \varphi_4 = 120^\circ$ and allowed for φ_2 and φ_3 the 120 and 125° values. For the four possible combinations of these φ_i values, the same θ_1, θ_2 combinations as before were considered and the 'total energies' for the model molecule in the crystal were calculated. The more interesting results are summarized in Fig. 6 and Table 5. In Table 5, the approximate geometry corresponding to the minimum calculated energy of the molecule in its crystalline environment is given.

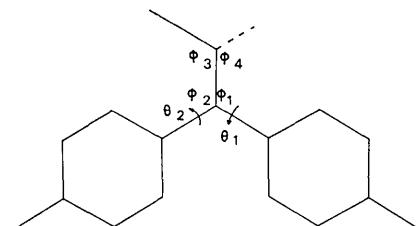


Fig. 5. Degrees of freedom allowed for the model molecule.

Table 5. Minimum energy geometries
in the model molecule after packing

Angles are in degrees and energy in Kcal.mole⁻¹ (see Fig. 6). For the most stable conditions (case A) the bond distances are C(1)-C(2)=1.35, C(2)-C(3)=1.48, C(2)-C(9)=1.53 Å.

θ_1	θ_2	φ_1	φ_2	φ_3	φ_4	Energy	Case
25	69	120	120	120	120	-673.13	A
25	69	120	120	125	120	-672.62	B
25	65	120	125	120	120	-673.08	C
25	65	120	125	125	120	-672.57	D

Discussion

The crystal structure of the molecule reveals the existence of three planes containing the two tolyl groups and the ethylenic group. The e.s.d.'s of the bond lengths and bond angles are probably underestimated, owing to the fact that the absorption correction was not allowed for. It seems that this effect is partially compensated by the values assumed by the thermal parameters. In any event, the thermal motion in the crystal appears to be very large, especially for the bromine atom and the two methyl groups. The carbon-bromine, the ethylenic and the two C_{sp^2} -C_{methyl} bond distances are in the expected range. The C(2)-C(3) and C(2)-C(9) bond lengths also show normal values; the difference between them, even if scarcely significant, is in

the correct sense, considering the different amount of rotation of the connected ring with respect to the ethylenic plane. The two dihedral angles are 24.4 and 67.9° (see Fig. 2).

The φ_2 and φ_3 angles are larger, at a 3σ level, than 120°, while φ_1 is not significantly different from 120°. The φ_4 angle has not been determined experimentally. No non-bonded distances significantly smaller than the corresponding 'touching distance' were found. On comparison with experimental values, the configuration of maximum stability for the isolated molecule (Table 4) shows a nice agreement of C-C bond lengths and φ_i angles; the θ_1 angle is also very close but the calculated θ_2 angle is significantly smaller than the observed one. A possible explanation of this discrepancy might be the interaction with neighbour molecules after packing in the crystal. This point was checked in the calculation of the energy of a crystal made of model molecules. The allowed degrees of freedom were chosen so that the energy needed for distortion is a minimum. The result is in good agreement with experiment (see Table 5); there is an increase of the calculated θ_2 angle to 69° with all the φ_i angles equal to 120° (case A), or to 65° with an increase of the φ_2 angle to 125° (case C). This suggest that the minimum lies between $\theta_2=65^\circ$ and 69° , with φ_2 between 120° and 125°.

The constancy of θ_1 in the isolated molecule and all the model crystals is particularly significant (see also Simonetta & Carrà, 1963). The agreement must not be overestimated owing to the many assumptions and approximations involved in the method of calculation and to the fact that small differences in energy are dealt with. There is no doubt, however, that the packing is responsible for the disagreement between the isolated and the real molecule geometries; we may have some confidence that the magnitude of the distortion due to packing is correct and that the geometry given for the isolated molecule can be considered calculated with a good approximation.

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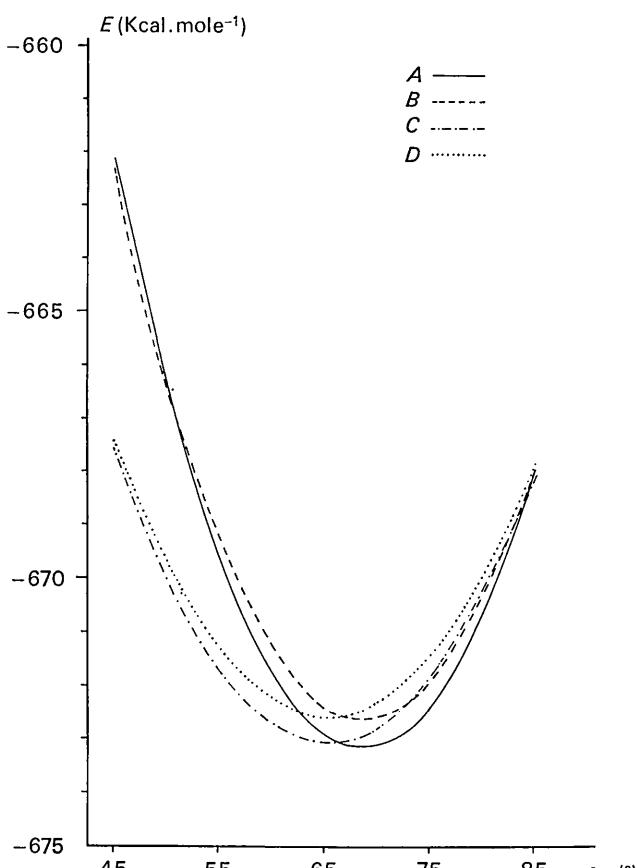


Fig. 6. 'Total energy' of the model molecule after packing (see Table 5).

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Structure Cristalline du Dichlorure de Cadmium Diformamide, $\text{CdCl}_2 \cdot 2\text{HCONH}_2$

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The crystal structure of bisformamide cadmium chloride has been determined and refined with three-dimensional data using Cu $K\alpha$ X-radiation. $\text{CdCl}_2 \cdot 2\text{HCONH}_2$ is triclinic with $a = 8.37 \pm 0.02$, $b = 7.26 \pm 0.02$, $c = 3.77 \pm 0.01$ Å, $\alpha = 93^\circ 41' \pm 20'$, $\beta = 109^\circ 07' \pm 20'$ and $\gamma = 113^\circ 48' \pm 20'$. $Z = 1$. Space group: $P\bar{1}$.

Fourier and least-squares methods were used to determine positional coordinates. The molecules of formamide are bonded to the metal by the oxygen atoms. The structure is built up from chains of octahedral CdCl_4O_2 polyhedra, linked along the [001] axis. The two O atoms belong to two molecules of formamide lying opposite to the plane containing the Cd atom and the four Cl atoms.

Introduction

Par la détermination des structures cristallines des composés d'addition moléculaire $\text{CdCl}_2 \cdot 2\text{OC}(\text{NH}_2)_2$ et $\text{CdCl}_2 \cdot 2\text{C}_2\text{H}_5\text{N}_3\text{O}_2$ Nardelli, Cavalca & Fava (1957) et Cavalca, Nardelli & Fava (1960) ont montré que la liaison Cd-X est généralement assurée par l'atome d'oxygène dans les composés du type $\text{CdCl}_2 \cdot 2\text{X}$, X désignant une molécule organique possédant des groupements C=O et C=N donneurs de paires d'électrons.

Le spectre d'absorption infra-rouge d'un composé d'addition moléculaire permet souvent de déterminer la nature de l'atome donneur lié au métal; l'existence d'une liaison métal-oxygène ou métal - azote entraîne généralement une diminution de la constante de force de la liaison C=O ou C=N et par là, une diminution de la fréquence $\nu(\text{C=O})$ ou $\nu(\text{C=N})$ par rapport à la molécule organique libre.

Pour $\text{CdCl}_2 \cdot 2\text{HCONH}_2$, le spectre d'absorption infra-rouge, enregistré en suspension dans le Nujol, ne permet pas de déterminer avec certitude si le formamide est lié au cadmium par l'atome d'oxygène ou par l'atome d'azote. En effet, la fréquence de la vibration d'elongation $\nu(\text{C-O})$ ne change pas par rapport à celle du formamide pur, à l'état liquide (Suzuki, 1960); la fréquence de la vibration $\nu(\text{C-N})$ augmente, au lieu de diminuer. Le Tableau 1 donne les fréquences des vibrations $\nu(\text{C-O})$ et $\nu(\text{C-N})$, du formamide lié au

Tableau 1. Fréquences des vibrations $\nu(\text{C-O})$ et $\nu(\text{C-N})$ du formamide engagé dans $\text{CdCl}_2 \cdot 2\text{HCONH}_2$ et du formamide non lié

Fréquences observées	$\text{CdCl}_2 \cdot 2\text{HCONH}_2$	HCONH_2
$\nu(\text{C-O})$	1690 cm^{-1}	1690 cm^{-1}
$\nu(\text{C-N})$	1350 cm^{-1}	1309 cm^{-1}